

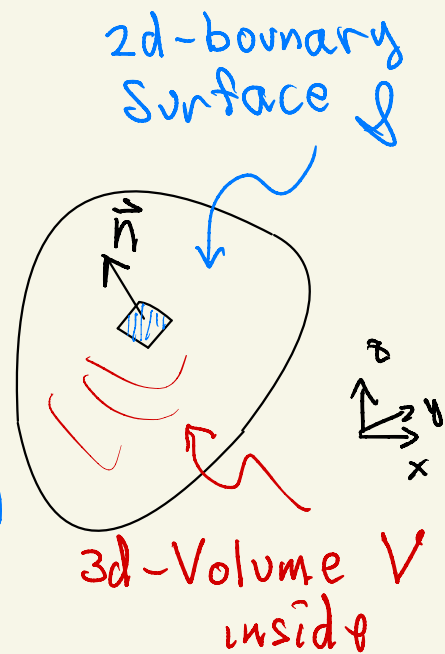
§ 16.4/5 Flux and Surface Integrals

①

Recall: the Divergence Theorem:

$$\iiint_V \text{Div } \vec{F} \, dV = \iint_{\partial} \vec{F} \cdot \vec{n} \, dS$$

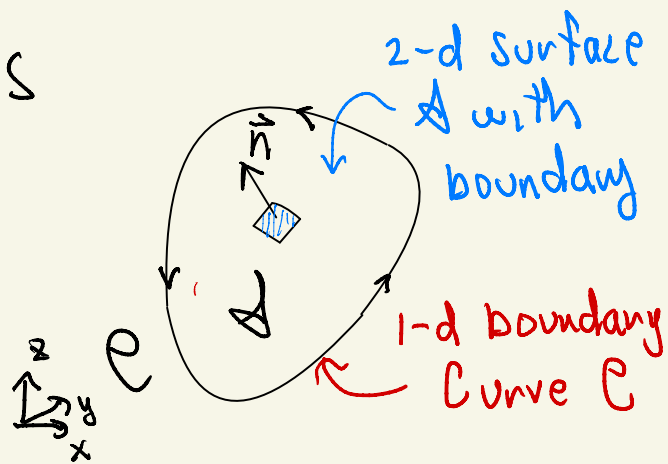
Flux of \vec{F} thru the closed boundary ∂



Recall that Stokes Theorem reads

$$\iint_{\partial} \text{Curl } \vec{F} \cdot \vec{n} \, dS = \int_C \vec{F} \cdot \vec{T} \, ds$$

Flux of $\text{Curl } \vec{F}$ thru ∂



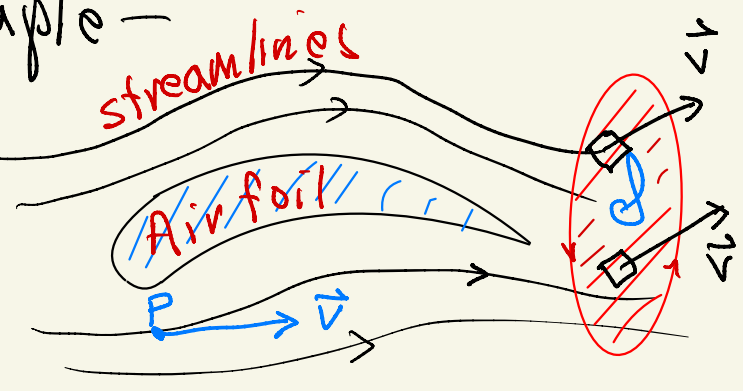
This reduces to Green's Thm in the special case when ∂ reduces to a region R in xy -plane inside see C in plane

Green's Thm: $\iint_R \underbrace{N_x - M_y}_{\text{Curl}(M, N, 0) \cdot (0, 0, 1)} \, dA = \int_C \vec{F} \cdot \vec{T} \, ds$

• Flux has a very important physical meaning:

Consider the fluid example -

A fluid particle P moves along curves $\vec{v}_P(t)$ called streamlines -



The fluid density $\rho(x, y, z) = \frac{\text{mass}}{\text{vol}}$ changes as the fluid "compresses" and "rarefies".

Q: Given a 2-d surface \mathcal{A} , what is the rate at which mass is passing thru \mathcal{A} ?

Ans: Flux = $\iint_{\mathcal{A}} \vec{F} \cdot \vec{n} dS = \frac{\text{Mass}}{\text{Time}}$ thru \mathcal{A}

$\vec{F} = \rho \vec{v}$ = mass flux vector.

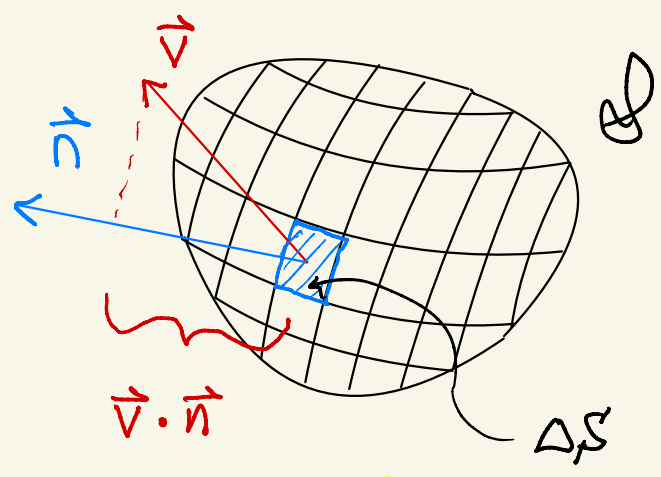
Check dimensions $[\rho \vec{v}] = \frac{\text{mass}}{\text{vol}} \frac{\text{dist}}{\text{time}} = \frac{M L}{L^3 T} = \frac{M}{L^2 T}$

mass / (area x time)

Conclude: $\rho \vec{v} = \frac{\text{mass}}{\text{area} \times \text{time}}$ moving

thru surface in direction \vec{v}

- Discretize \mathcal{S} into a grid of small areas ΔS



- Only the component of $\vec{v} \perp \Delta S$ contributes

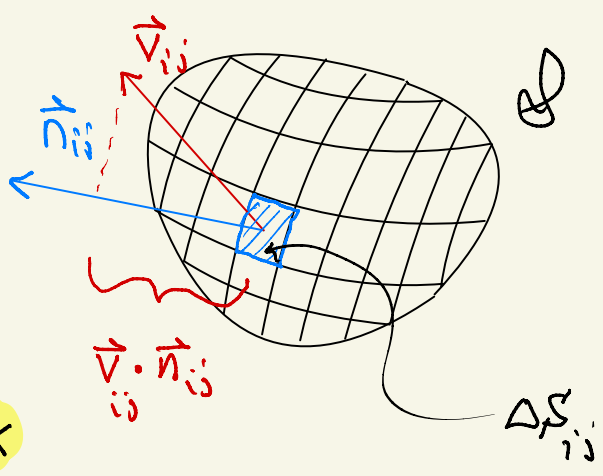
to mass flow thru ΔS . This is $\vec{v} \cdot \vec{n}$

- $\delta \vec{v} \cdot \vec{n} = \frac{\text{mass}}{\text{vol}} \frac{\text{dist}}{\text{time}}$ moving \perp to ΔS

$= \frac{\text{mass}}{\text{area time}}$ out thru ΔS

- $\delta \vec{v} \cdot \vec{n} \Delta S = \frac{\text{mass}}{\text{time}}$ out thru ΔS

- $\frac{\text{Total Mass}}{\text{Time}} \approx \sum_{ij} \delta_{ij} \vec{v}_{ij} \cdot \vec{n}_{ij} \Delta S_{ij}$
out thru \mathcal{S}



$$\frac{\text{Total Mass}}{\text{Time}} = \lim_{N \rightarrow \infty} \sum_{ij} \delta_{ij} \vec{v}_{ij} \cdot \vec{n}_{ij} \Delta S_{ij}$$

$$= \iint_{\mathcal{S}} \delta \vec{v} \cdot \vec{n} dS = \text{Flux}$$

• The argument works for any density - (4)

$$\delta = \frac{\text{charge}}{\text{vol}}, \frac{\text{contaminant}}{\text{vol}}, \frac{\text{"stuff"}}{\text{vol}}$$

- assume its moving at velocity \vec{v}

- define $\delta \vec{v} = \text{"stuff" flux vector}$

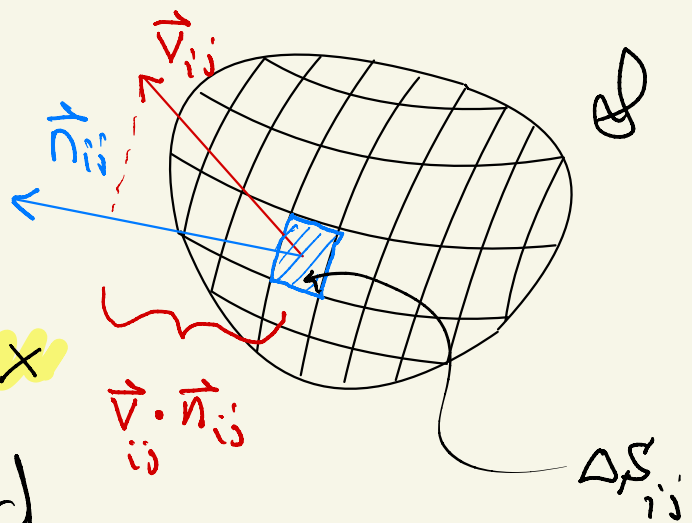
- then the "Flux thru \mathcal{S} " is defined as

$$\iint_{\mathcal{S}} \delta \vec{v} \cdot \vec{n} \, d\mathcal{S} = \text{"rate at which stuff is passing thru } \mathcal{S} \text{ in direction of normal } \vec{n} \text{"}$$

Precisely -

$$\frac{\text{Total Mass}}{\text{Time}} = \lim_{N \rightarrow \infty} \sum_{ij} \delta_{ij} \vec{v}_{ij} \cdot \vec{n}_{ij} \Delta \mathcal{S}_{ij}$$

$$= \iint_{\mathcal{S}} \delta \vec{v} \cdot \vec{n} \, d\mathcal{S} = \text{Flux}$$



To calculate this we need

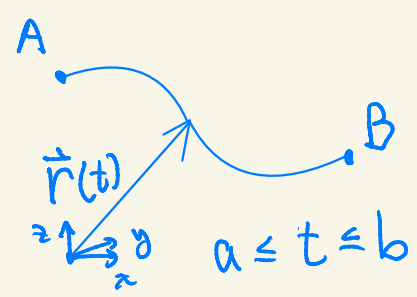
to define the Riemann Sum in a Coordinate System \mathcal{D}

Calculating surface integrals -

Turns out - there is a very natural way to define Flux integrals in terms of a coordinate system on the surface - analogous to coordinate systems on curves -

Defn: we call a coord system a parameterization

Recall: $\int_C \vec{F} \cdot \vec{T} ds = \int_a^b \vec{F} \cdot \vec{v} dt$



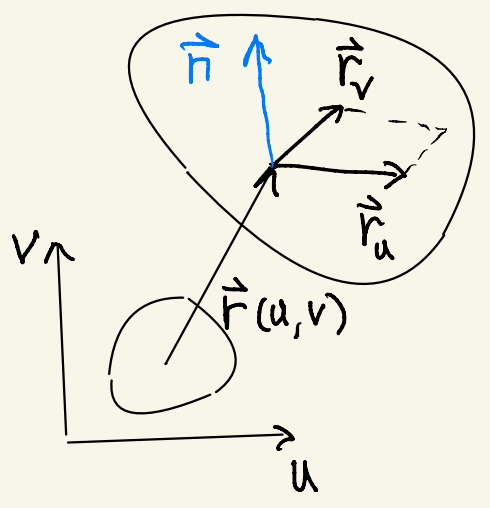
Similarly for surfaces -

$$\vec{r}(u,v) = (x(u,v), y(u,v), z(u,v))$$

$$\vec{r}_u = (x_u(u,v), y_u(u,v), z_u(u,v))$$

$$\vec{r}_v = (x_v(u,v), y_v(u,v), z_v(u,v))$$

$$\vec{n} = \frac{\vec{r}_u \times \vec{r}_v}{\|\vec{r}_u \times \vec{r}_v\|} = \text{unit normal}$$



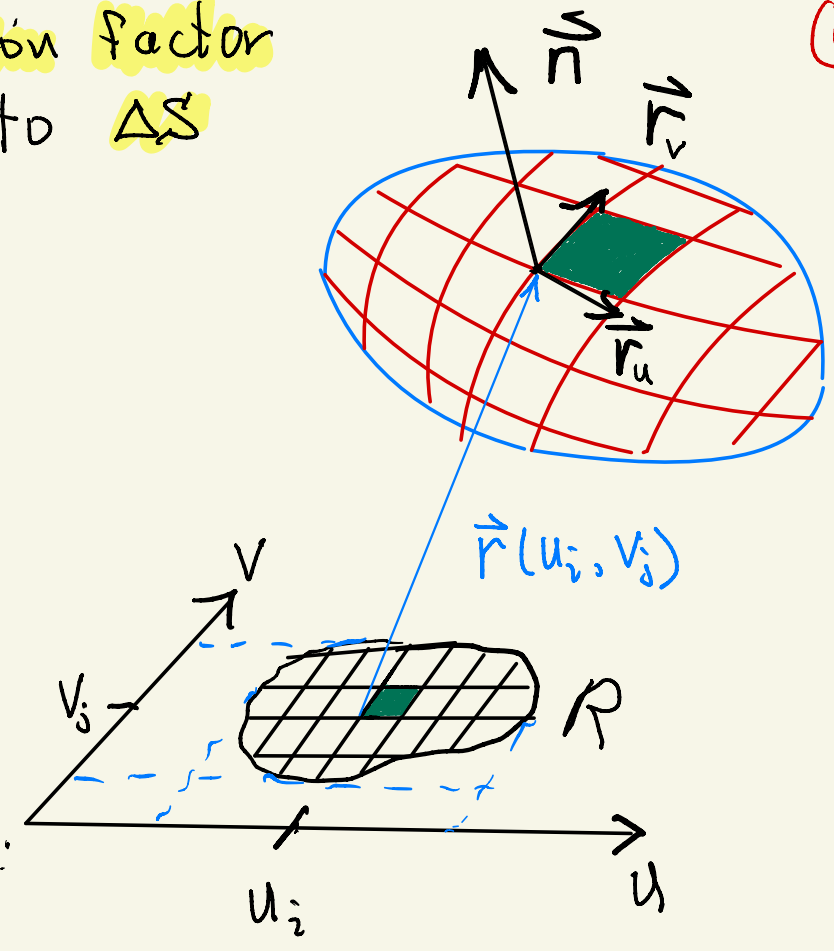
It remains to get the amplification factor -

Turns out: The amplification factor A that scales $\Delta u \Delta v$ to ΔS

$$\Delta S = A \Delta u \Delta v$$

is $A = |\vec{r}_u \times \vec{r}_v|$

From this fact we can compute a surface integral flux in a coordinate system.



$$\iint_S \vec{F} \cdot \vec{n} \, dS = \lim_{N \rightarrow \infty} \sum \underbrace{\vec{F}_{ij} \cdot \vec{n}_{ij}}_{\frac{\vec{r}_u \times \vec{r}_v}{\|\vec{r}_u \times \vec{r}_v\|}} \underbrace{\Delta S_{ij}}_{\|\vec{r}_u \times \vec{r}_v\| \Delta u \Delta v}$$

$$= \lim_{N \rightarrow \infty} \sum_{(u_i, v_i) \in R} \underbrace{\vec{F}_{ij} \cdot \frac{\vec{r}_u \times \vec{r}_v}{\|\vec{r}_u \times \vec{r}_v\|} \|\vec{r}_u \times \vec{r}_v\| \Delta u \Delta v}_{\text{Riemann Sum in } (u, v)}$$

$$= \iint_R \vec{F} \cdot (\vec{r}_u \times \vec{r}_v) \, dA$$

A Chapter 15 integral over a region R in uv -plane

Summary →

Line Integral: $\int_C \vec{F} \cdot \vec{T} \, ds = \int_a^b \vec{F} \cdot \vec{v} \, dt$

$\vec{r}(t): a \leq t \leq b$

Surface Integral: $\iint_S \vec{F} \cdot \vec{n} \, dS = \iint_R \vec{F} \cdot (\vec{r}_u \times \vec{r}_v) \, dA$

$\vec{r}(u, v): (u, v) \in R$

The only step that needs justification is the amplification factor -

Q: why is $dS = |\vec{r}_u \times \vec{r}_v| \, du \, dv$?

We first recall the cross-product.

Cross Product: $\vec{A} = (a_1, a_2, a_3)$ $\vec{B} = (b_1, b_2, b_3)$

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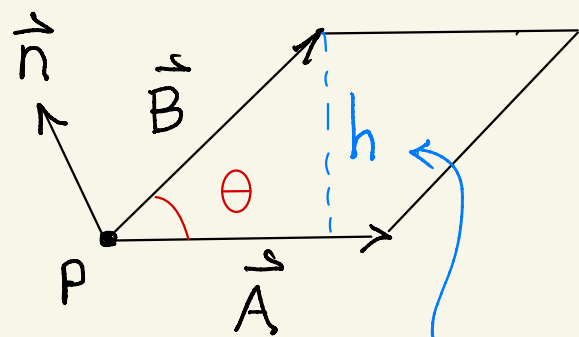
How to calculate it:

$$\vec{A} \times \vec{B} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix} = (a_2 b_3 - a_3 b_2) \hat{i} - (a_1 b_3 - a_3 b_2) \hat{j} + (a_1 b_2 - a_2 b_1) \hat{k}$$

Geometrical meaning:

$$\vec{A} \times \vec{B} = \underbrace{\|\vec{A}\| \|\vec{B}\| \sin \theta}_{h} \hat{n}$$

Area of the
parallelogram



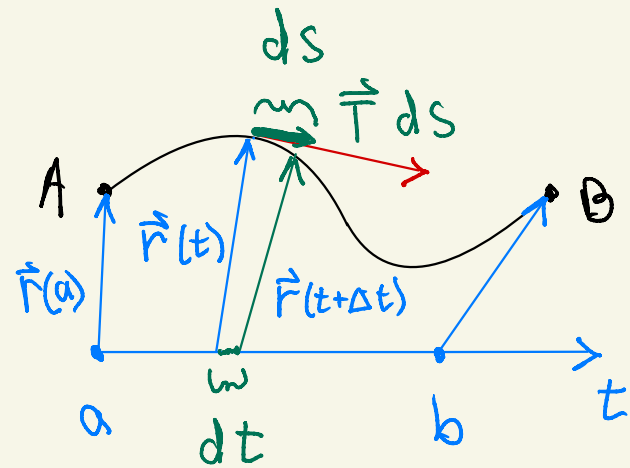
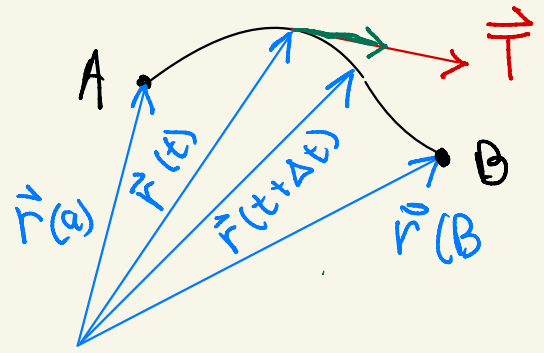
parallelogram
determined by
 \vec{A} and \vec{B}

Conclude: The cross product points in
direction \perp \vec{A} & \vec{B} (direction by right hand rule)
and has a length = area of parallelogram.

Now consider a curve \mathcal{C} :

- At a point $\vec{r}(t)$ on a curve, the vector $\vec{T} ds$ points tangent to \mathcal{C} , and has a length $ds = |\vec{r}'(t)| dt$

(ds is distance along tangent line, a good approximation to Δs along curve, when ds is small)



- Similarly: Given a surface $\vec{r}(u, v)$

$\vec{r}_u du = \vec{T}_u ds$ is the vector on the side of Δs tangent to curve $\vec{r}_v(u)$

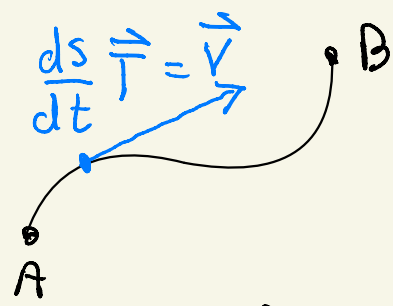
$\vec{r}_v dv = \vec{T}_v ds$ is the vector on the side of Δs tangent to curve $\vec{r}_u(v)$

Explanation: Why $|\vec{r}_u \times \vec{r}_v|$ is the amplification factor for area I.e.,

$$dS = |\vec{r}_u \times \vec{r}_v| du dv$$

(1) First recall that for a curve $\vec{r}(t)$, we have $\|\vec{v}\| = \|\vec{r}'(t)\| = \frac{ds}{dt} \Rightarrow ds = \|\vec{r}'(t)\| dt$

thus $\vec{T} ds$ is a vector of length ds pointing tangent



(2) Similarly, $\vec{r}(u, v)$ parameterizes a surface - and at fixed v

$\vec{r}(u, v) = \vec{r}_v(u)$ is a curve with parameter u

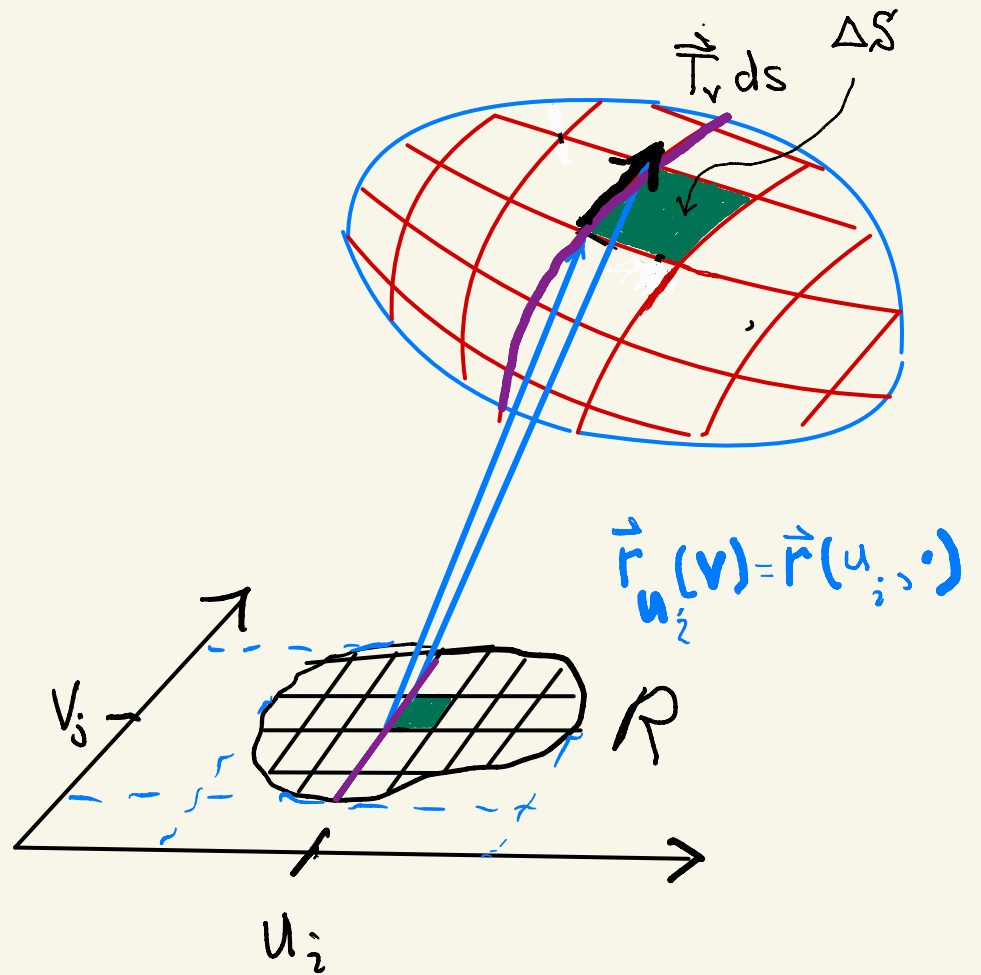
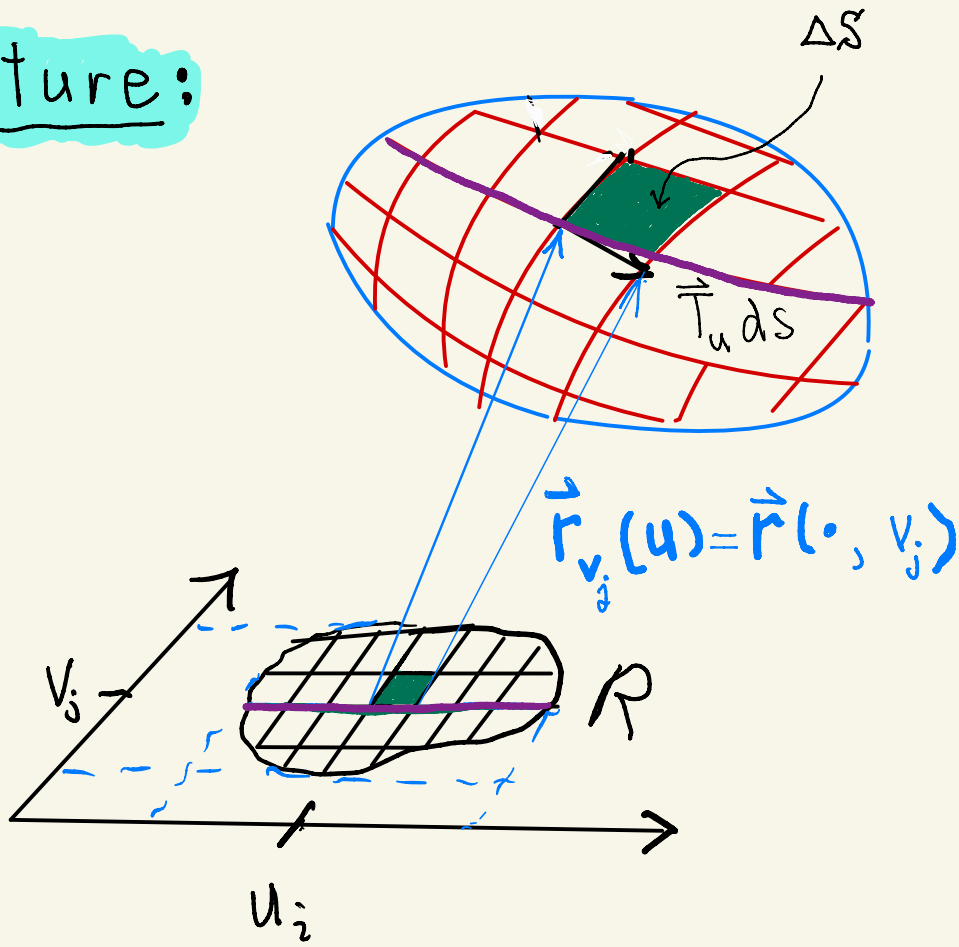
$$\frac{\partial \vec{r}}{\partial u}(u, v) = \frac{d\vec{r}_v}{du} = \vec{T}_u ds \text{ is one side of parallelogram } \Delta S$$

Same for v :

$\vec{r}(u, v) = \vec{r}_u(v)$ is a curve with parameter v

$$\frac{\partial \vec{r}}{\partial v}(u, v) = \frac{d\vec{r}_u}{dv} = \vec{T}_v ds \text{ is the other side of parallelogram } \Delta S$$

Picture:

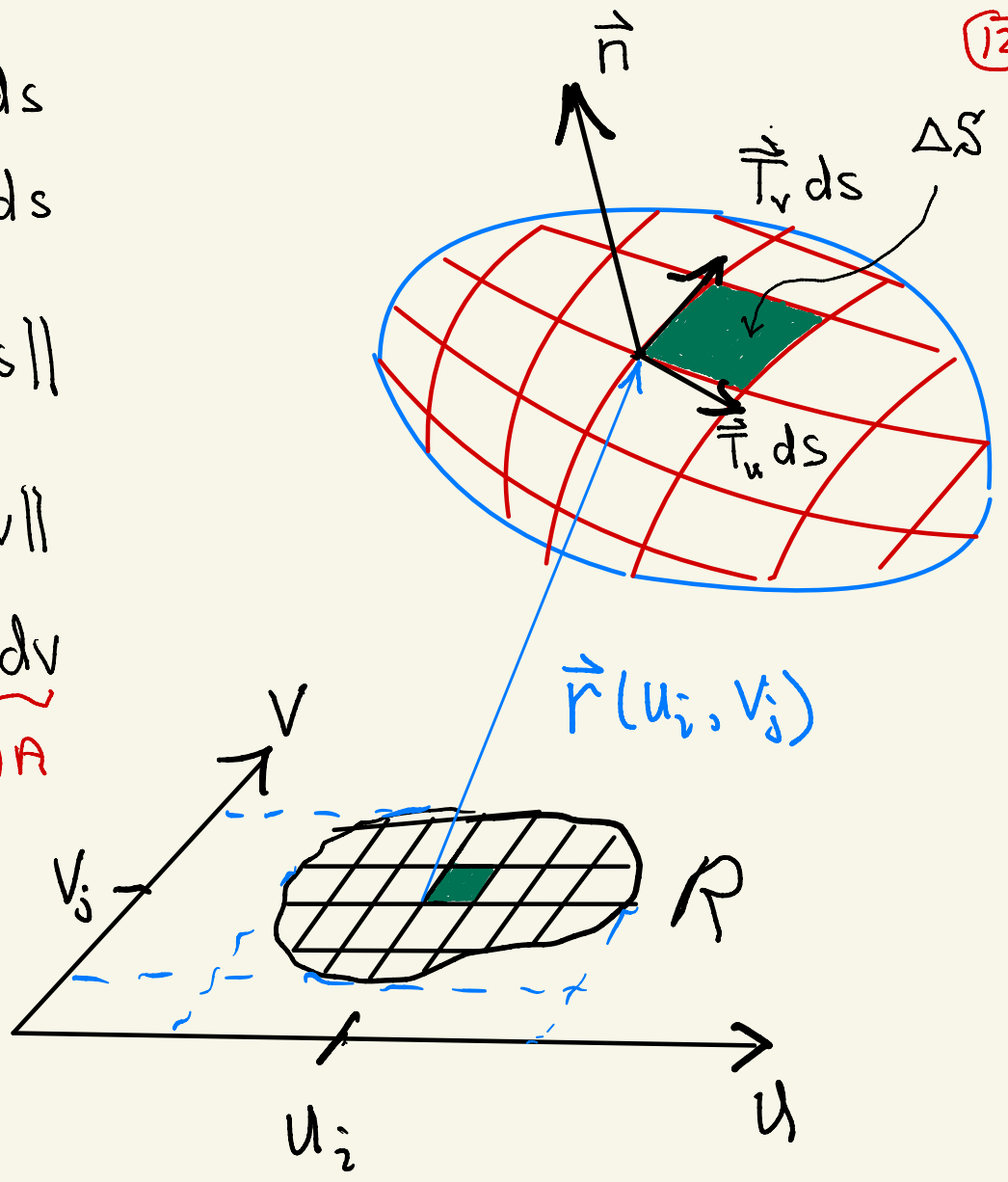


(3) $\vec{r}_u du = \vec{T}_u ds$
 $\vec{r}_v dv = \vec{T}_v ds$

$$dS = \|\vec{T}_u ds \times \vec{T}_v ds\|$$

$$= \|\vec{r}_u du \times \vec{r}_v dv\|$$

$$= \underbrace{\|\vec{r}_u \times \vec{r}_v\|}_{\text{Amplification Factor for Area}} \underbrace{du dv}_{dA}$$



Summary:

$$dS = \|\vec{r}_u \times \vec{r}_v\| dA$$

so

$$\iint_{R_{uv}} \vec{F} \cdot \vec{n} dS = \iint_{R_{uv}} \vec{F} \cdot \underbrace{\frac{\vec{r}_u \times \vec{r}_v}{\|\vec{r}_u \times \vec{r}_v\|}}_{\vec{n}} \underbrace{\|\vec{r}_u \times \vec{r}_v\|}_{dS} dA$$

$$= \iint_{R_{uv}} \vec{F} \cdot \vec{r}_u \times \vec{r}_v dA$$

Example: Assume a surface \mathcal{S} is given by

$$\vec{r}(u, v) = \overrightarrow{(u-v, u+v, u+2v)} \quad \begin{array}{l} 0 \leq u \leq 1 \\ 0 \leq v \leq 2 \end{array}$$

$\underbrace{\hspace{1.5cm}}_{x(u,v)} \quad \underbrace{\hspace{1.5cm}}_{y(u,v)} \quad \underbrace{\hspace{1.5cm}}_{z(u,v)}$

Assume a density $\delta(x, y, z) = x \frac{\text{kg}}{\text{m}^3}$ is moving thru the surface at velocity $\vec{v} = y \overrightarrow{(1, 2, 1)} \frac{\text{m}}{\text{s}}$

Find the rate and direction at which mass is passing thru the surface.

Soln: $\vec{F} = \delta \vec{v} = xy \overrightarrow{(1, 2, 1)}$ is the mass flux vector

Flux = $\iint_{\mathcal{S}} \vec{F} \cdot \vec{n} \, ds = \frac{\text{mass}}{\text{time}}$ thru \mathcal{S} in direction of normal vector \vec{n} .

$$\vec{r}_u = \frac{\partial \vec{r}}{\partial u} = \frac{\partial}{\partial u} \overrightarrow{(u-v, u+v, u+2v)} = \overrightarrow{(1, 1, 1)}$$

$$\vec{r}_v = \frac{\partial \vec{r}}{\partial v} = \frac{\partial}{\partial v} \overrightarrow{(u-v, u+v, u+2v)} = \overrightarrow{(-1, 1, 2)}$$

$$\vec{r}_3 \times \vec{r}_5 = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ 2 & 1 & 1 \\ -1 & 1 & 2 \end{vmatrix} = \hat{i}(2-1) - \hat{j}(2-(-1)) + \hat{k}(1-(-1))$$

$$= \hat{i} - 3\hat{j} + 2\hat{k} = \overrightarrow{(1, -3, 2)}$$

$$\vec{n} = \frac{\vec{r}_u \times \vec{r}_v}{\|\vec{r}_u \times \vec{r}_v\|} = \frac{\overrightarrow{(1, -3, 2)}}{\sqrt{1^2 + (-3)^2 + 2^2}} = \frac{1}{\sqrt{14}} \overrightarrow{(1, -3, 2)}$$

↑
positive

Since z-component of $\vec{n} > 0$, this is upward normal

$$\iint_S \vec{F} \cdot \vec{n} \, dS = \iint_{R_{uv}} \underbrace{xy}_{\vec{F}} \underbrace{\frac{1}{\sqrt{14}} \overrightarrow{(1, -3, 2)}}_{\vec{n}} \underbrace{\sqrt{14} \, dA}_{dS}$$

$x = u - v, y = u + v$

$$= \int_0^2 \int_0^1 (u-v)(u+v)(1-6+2) \, du \, dv = \iint_{R_{uv}} \vec{F} \cdot \vec{r}_u \times \vec{r}_v \, dA$$

$$= -3 \int_0^2 \int_0^1 u^2 - v^2 \, du \, dv = -3 \int_0^2 \left[\frac{u^3}{3} - v^2 u \right]_{u=0}^{u=1} dv = -3 \int_0^2 \left[\frac{1}{3} - v^2 \right] dv$$

$$= -3 \left[\frac{1}{3}v - \frac{v^3}{3} \right]_{v=0}^{v=2} = -3 \left[\left(\frac{2}{3} - \frac{8}{3} \right) - 0 \right] = 6 \frac{kg}{s} > 0$$

⇒ upward thru \mathcal{S} ✓